

## Lecture 10:

### Variational Methods I: Basic Ideas

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#### Why Variational Methods?

### Why Variational Methods?

- ◆ image transformation that satisfies an optimality criterion
- ◆ clear description of all model assumptions without any hidden constraints
- ◆ want to find an image restoration model where
  - the filtered image  $u$  is close to the original image  $f$
  - the filtered image  $u$  should be (at least piecewise) smooth
- ◆ treatment at the boundaries should result in a natural way from the model assumptions

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## Modelling

### Basic Idea

(Bertero et al. 1988, Nordström 1990, Schnörr 1994, Charbonnier et al. 1994)

- ◆ variational restoration of some image  $f$ : minimiser  $u_\alpha$  of

$$E_f(u) := \int_{\Omega} \left( \underbrace{(u-f)^2}_{\text{similarity}} + \alpha \underbrace{\Psi(|\nabla u|^2)}_{\text{smoothness}} \right) dx$$

- ◆ first term (*data term, similarity term, fidelity term*): rewards similarity to the original image
- ◆ second term (*smoothness term, regulariser, penaliser*): penalises deviations from (piecewise) smoothness
- ◆  $\alpha > 0$ : *regularisation parameter (smoothness weight)*
- ◆ Variational methods of this type are also called *regularisation methods*.

### Technical Assumptions on the Penalising Function $\Psi$ :

- ◆  $\Psi$  is differentiable and increasing:  $\Psi'(s^2) > 0$ .
- ◆  $\Psi(s^2)$  is convex in  $s$ .
- ◆ There exist constants  $c_1, c_2 > 0$  such that  $c_1 s^2 \leq \Psi(s^2) \leq c_2 s^2$  for all  $s^2$ .

### Examples:

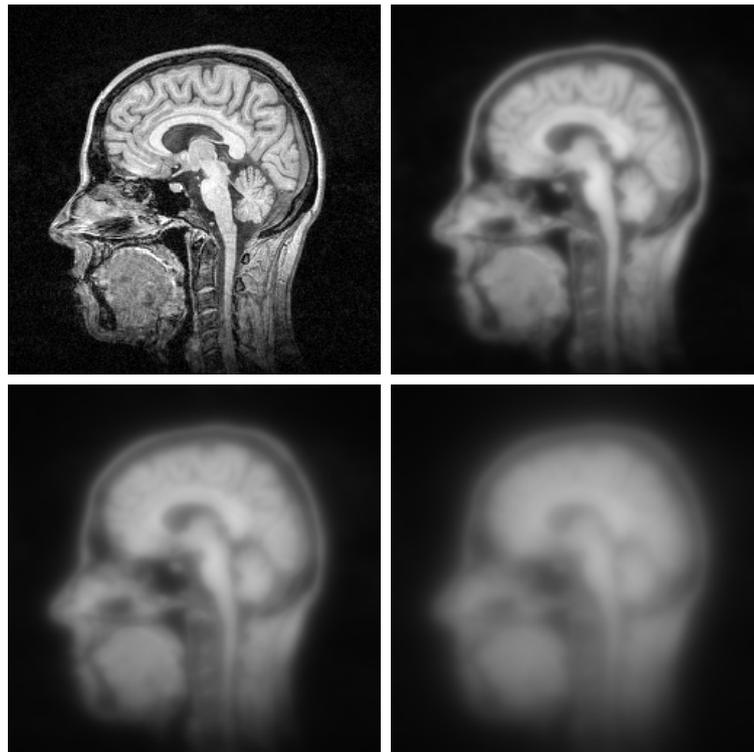
- ◆ Whittaker 1923, Tikhonov 1963:

$$\Psi(s^2) = s^2$$

- ◆ Charbonnier et al. 1994:

$$\Psi(s^2) = 2\lambda^2 \sqrt{1 + s^2/\lambda^2} - 2\lambda^2$$

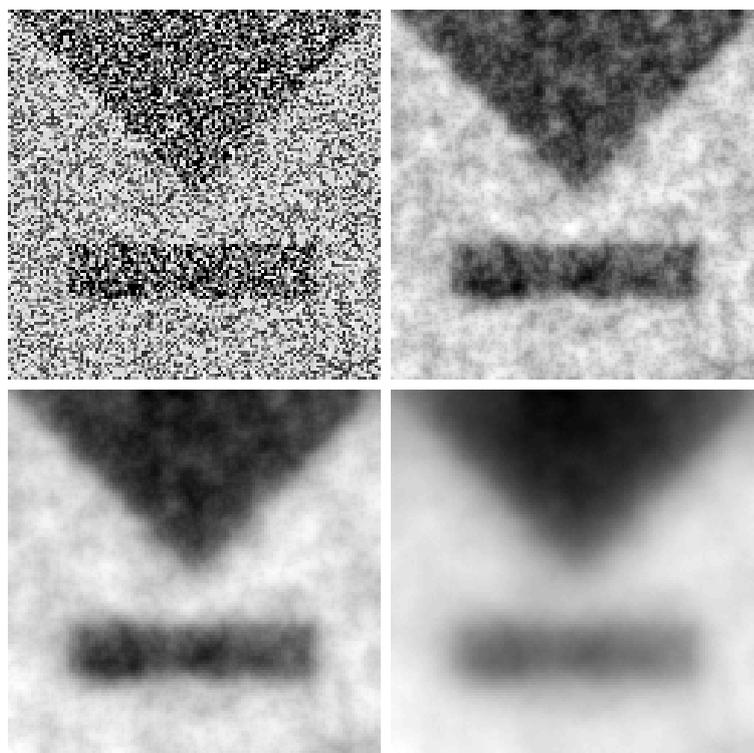
### Modelling (3)



Variational image restoration with the Whittaker–Tikhonov penaliser. **Top left to bottom right:** Regularisation parameters  $\alpha = 0, 10, 30, 100$ .

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### Modelling (4)

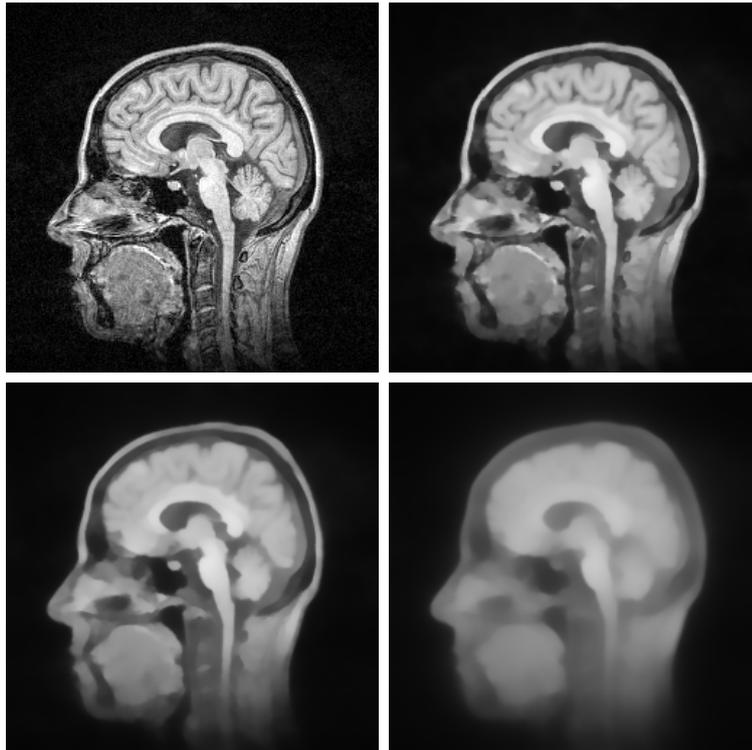


Variational image restoration with the Whittaker–Tikhonov penaliser. **Top left to bottom right:** Regularisation parameters  $\alpha = 0, 5, 20, 100$ .

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## Modelling (5)

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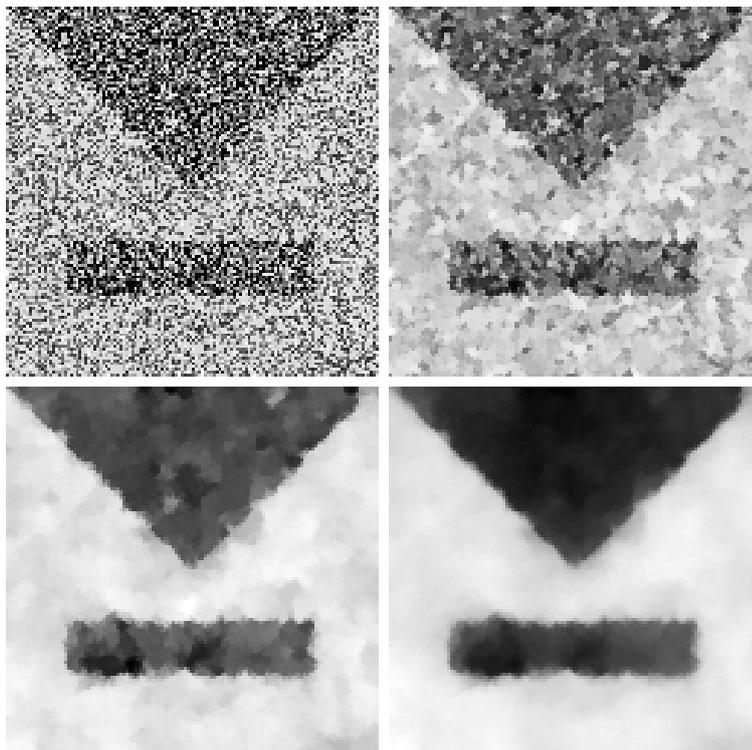


Variational image restoration with the Charbonnier penaliser ( $\lambda = 2$ ). **Top left to bottom right:** Regularisation parameters  $\alpha = 0, 10, 30, 100$ .

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## Modelling (6)

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Variational image restoration with the Charbonnier penaliser ( $\lambda = 1$ ). **Top left to bottom right:** Regularisation parameters  $\alpha = 0, 20, 50, 100$ .

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## Well-Posedness and Scale-Space Properties

It seems that the regularisation parameter  $\alpha$  plays the role of a scale parameter. The following well-posedness and scale-space properties can be established:

(a) **Well-Posedness and Regularity** (Schnörr 1994)

- Existence of a unique minimizer  $u_\alpha$  for bounded image  $f$ .
- Smoothness:  $u_\alpha \in H^2(\Omega)$  (“twice differentiable”)
- $u_\alpha$  depends continuously on  $f$ .
- No contrast enhancement for convex regularisers!

(b) **Average Grey Level Invariance**

$$\frac{1}{|\Omega|} \int_{\Omega} u_\alpha(\mathbf{x}) \, d\mathbf{x} = \frac{1}{|\Omega|} \int_{\Omega} f(\mathbf{x}) \, d\mathbf{x} =: \mu.$$

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(c) **Maximum–Minimum Principle**

$$\inf f \leq u_\alpha(\mathbf{x}) \leq \sup f \quad \forall \mathbf{x} \in \Omega, \forall \alpha > 0$$

(d) **Lyapunov Property** (Scherzer/W. 2000)

Let  $r \in C^2$  be convex. Then

$$V(\alpha) := \int_{\Omega} r(u_\alpha(\mathbf{x})) \, d\mathbf{x}$$

is bounded from below and satisfies  $V(\alpha) \leq V(0)$ .

(e) **Convergence** (Scherzer/W. 2000)

$u_\alpha(\mathbf{x})$  converges to the average grey level  $\mu$  for  $\alpha \rightarrow \infty$   
(in 2D: convergence in  $L^p(\Omega)$ ,  $1 \leq p < \infty$ ).

These theoretical results resemble the ones for diffusion filtering (Lectures 4 and 8). Is there a specific reason behind this similarity?

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## Euler–Lagrange Equations

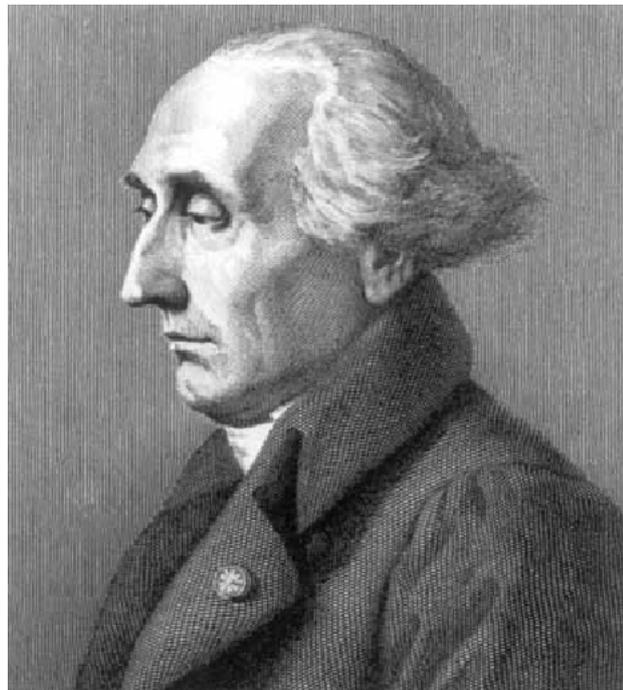
## Standard Calculus:

- ◆ considers real-valued *functions*  $f(x)$ , i.e. mappings from a *number*  $x$  into  $\mathbb{R}$ .
- ◆ If  $f$  has a minimum in  $\xi$ , then it necessarily holds  $f'(\xi) = 0$ .
- ◆ If  $f$  is strictly convex and  $f'(\xi) = 0$ , then  $\xi$  is the unique minimum of  $f$ .

## Calculus of Variations:

- ◆ considers real-valued *functionals*  $E(u)$ , i.e. mappings from a *function*  $u(x)$  into  $\mathbb{R}$ .
- ◆ If  $E$  is minimised by a function  $v$ , then  $v$  has to satisfy necessarily a so-called *Euler–Lagrange equation*. This is a partial differential equation in  $v$ .
- ◆ If the  $E$  is strictly convex and satisfies the Euler–Lagrange equation, then  $v$  is the unique minimiser of  $E$ .

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The mathematicians Euler and Lagrange belong to the founders of the calculus of variations. **Left:** Leonhard Euler (1707–1783). **Right:** Joseph–Louis Lagrange (1736–1813).

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## Euler–Lagrange Equations (3)

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### Theorem (Euler-Lagrange Equations in 1-D)

A smooth function  $u(x)$ ,  $x \in [a, b]$  that minimises the 1-D energy functional

$$E(u) = \int_a^b F(x, u, u') dx$$

satisfies necessarily the *Euler–Lagrange equation*

$$F_u - \frac{d}{dx} F_{u'} = 0$$

and the so-called *natural boundary conditions*

$$F_{u'} = 0$$

for  $x = a$  and  $x = b$ .

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## Euler–Lagrange Equations (4)

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### Proof

We assume that  $v(x)$  is a sufficiently often differentiable minimiser of  $E$ . We embed  $v(x)$  into the family

$$u(x, \varepsilon) := v(x) + \varepsilon h(x)$$

with some perturbation function  $h(x)$ .

Since  $v(x)$  minimises  $E(u)$ , we know that the scalar-valued function

$$g(\varepsilon) := E(u(x, \varepsilon)) = E(v + \varepsilon h)$$

has a minimum in  $\varepsilon = 0$ . Therefore, we have

$$0 = g'(0) = \frac{d}{d\varepsilon} E(v + \varepsilon h) \Big|_{\varepsilon=0} = \frac{d}{d\varepsilon} \int_a^b F(x, \underbrace{v + \varepsilon h}_{u(\cdot, \varepsilon)}, \underbrace{v' + \varepsilon h'}_{u'(\cdot, \varepsilon)}) dx \Big|_{\varepsilon=0}.$$

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## Euler–Lagrange Equations (5)

By means of the chain rule it follows that

$$\begin{aligned} 0 &= \int_a^b \left( F_u(x, u, u') h(x) + F_{u'}(x, u, u') h'(x) \right) dx \Big|_{\varepsilon=0} \\ &= \int_a^b \left( F_u(x, v, v') h(x) + F_{u'}(x, v, v') h'(x) \right) dx. \end{aligned}$$

Partial integration of the second term yields

$$0 = \int_a^b \left( F_u(x, v, v') h(x) - \frac{d}{dx} F_{u'}(x, v, v') h(x) \right) dx + F_{u'}(x, v, v') h(x) \Big|_{x=a}^{x=b} \quad (1)$$

If  $v$  is a minimiser within some family of competing functions  $u(\varepsilon, x) = v(x) + \varepsilon h(x)$ , then  $v$  is also minimiser within the smaller class of functions where the perturbation  $h(x)$  satisfies  $h(a) = 0 = h(b)$ . Thus,

$$0 = \int_a^b \left( F_u(x, v, v') - \frac{d}{dx} F_{u'}(x, v, v') \right) h(x) dx.$$

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## Euler–Lagrange Equations (6)

This gives the Euler–Lagrange equation

$$F_u - \frac{d}{dx} F_{u'} = 0.$$

Note that this equation does not depend on the perturbation  $h$ .

Plugging it into Equation (1) gives

$$F_{u'}(x, v, v') h(x) \Big|_{x=a}^{x=b} = 0.$$

which holds for arbitrary perturbations  $h$  (also with  $h(a) \neq 0$  and  $h(b) \neq 0$ ).

Thus, one obtains the natural boundary conditions

$$F_{u'} = 0$$

for  $x = a$  und  $x = b$ . This concludes the proof.

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## Euler–Lagrange Equations (7)



### Theorem (Euler-Lagrange Equations in 2-D)

A minimiser of the 2-D energy functional

$$E(u) := \int_{\Omega} F(x_1, x_2, u, u_{x_1}, u_{x_2}) dx$$

satisfies necessarily the Euler–Lagrange equation

$$F_u - \partial_{x_1} F_{u_{x_1}} - \partial_{x_2} F_{u_{x_2}} = 0$$

with the natural boundary conditions

$$\mathbf{n}^\top \begin{pmatrix} F_{u_{x_1}} \\ F_{u_{x_2}} \end{pmatrix} = 0$$

at the image boundary  $\partial\Omega$  with normal vector  $\mathbf{n}$ .

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## Euler–Lagrange Equations (8)



### Euler-Lagrange Equations Applied to our Functional

For our energy functional

$$E_f(u) := \int_{\Omega} \left( (u-f)^2 + \alpha \Psi(|\nabla u|^2) \right) dx$$

we have

$$\begin{aligned} F &= (u-f)^2 + \alpha \Psi(u_{x_1}^2 + u_{x_2}^2), \\ F_u &= 2(u-f), \\ F_{u_{x_1}} &= 2\alpha u_{x_1} \Psi'(u_{x_1}^2 + u_{x_2}^2), \\ F_{u_{x_2}} &= 2\alpha u_{x_2} \Psi'(u_{x_1}^2 + u_{x_2}^2). \end{aligned}$$

All we have to do now is to plug this into the Euler-Lagrange equation and the natural boundary conditions.

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## Euler–Lagrange Equations (9)



The **Euler-Lagrange equation**  $F_u - \partial_{x_1} F_{u_{x_1}} - \partial_{x_2} F_{u_{x_2}} = 0$  becomes

$$2(u - f) - 2\alpha \partial_{x_1} \left( \Psi'(u_{x_1}^2 + u_{x_2}^2) u_{x_1} \right) - 2\alpha \partial_{x_2} \left( \Psi'(u_{x_1}^2 + u_{x_2}^2) u_{x_2} \right) = 0.$$

Division by 2 yields

$$u - f - \alpha \operatorname{div} (\Psi'(|\nabla u|^2) \nabla u) = 0.$$

The **boundary condition**  $\mathbf{n}^\top \begin{pmatrix} F_{u_{x_1}} \\ F_{u_{x_2}} \end{pmatrix} = 0$  becomes

$$0 = \mathbf{n}^\top \begin{pmatrix} 2\alpha u_{x_1} \Psi'(u_{x_1}^2 + u_{x_2}^2) \\ 2\alpha u_{x_2} \Psi'(u_{x_1}^2 + u_{x_2}^2) \end{pmatrix} = 2\alpha \Psi'(|\nabla u|^2) \mathbf{n}^\top \nabla u.$$

Since  $\alpha > 0$  and  $\Psi'(|\nabla u|^2) > 0$ , we obtain the Neumann boundary condition

$$0 = \mathbf{n}^\top \nabla u = \partial_{\mathbf{n}} u.$$

## Relations to Diffusion Filters (1)



### Relations to Diffusion Filters (Scherzer/W. 2000)

Let us rewrite the Euler–Lagrange equation as

$$\frac{u - f}{\alpha} = \operatorname{div} (\Psi'(|\nabla u|^2) \nabla u).$$

This can be regarded as fully implicit time discretisation of the diffusion filter

$$\begin{aligned} \partial_t u &= \operatorname{div} (\Psi'(|\nabla u|^2) \nabla u), \\ u(\mathbf{x}, 0) &= f(\mathbf{x}) \end{aligned}$$

with a single time step of size  $\alpha$ .

In both cases we have Neumann boundary conditions.

*A variational method with regulariser  $\Psi(|\nabla u|^2)$  and regularisation parameter  $\alpha$  approximates a diffusion filter with diffusivity  $\Psi'(|\nabla u|^2)$  and stopping time  $\alpha$ .*

We may thus regard the well-posedness and scale-space theory for variational methods as a time-discrete theory for diffusion filters. This completes the continuous, space-discrete and fully discrete diffusion framework of Lecture 8.

Regularisers and Their Corresponding Diffusivities

- ◆ The **Tikhonov regulariser**

$$\Psi(|\nabla u|^2) = |\nabla u|^2$$

corresponds to the linear diffusivity

$$\Psi'(|\nabla u|^2) = 1.$$

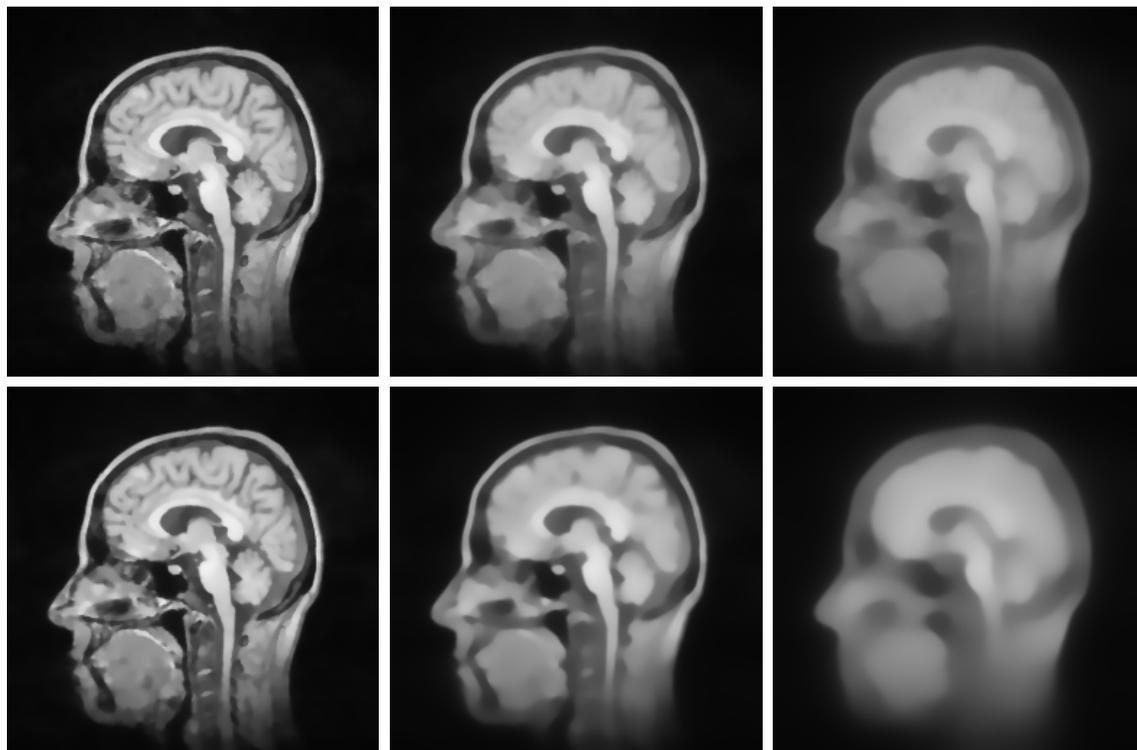
- ◆ The **Charbonnier regulariser**

$$\Psi(|\nabla u|^2) = 2\lambda^2\sqrt{1 + |\nabla u|^2/\lambda^2} - 2\lambda^2$$

leads to the nonlinear Charbonnier diffusivity

$$\Psi'(|\nabla u|^2) = \frac{1}{\sqrt{1 + |\nabla u|^2/\lambda^2}}.$$

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Comparison between variational image restoration and nonlinear diffusion filtering (Charbonnier diffusivity,  $\lambda = 2$ ). (b) **Top:** Variational restoration with regularisation parameters  $\alpha = 10, 30, 100$ . (a) **Bottom:** Temporal evolution of the diffusion filter,  $t = 10, 30, 100$ .

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## Summary

- ◆ Variational image restoration methods minimise an energy functional with a similarity term and a smoothness term.
- ◆ They lead to a similar well-posedness and scale-space theory as nonlinear diffusion filters.
- ◆ The Euler-Lagrange equation provides a necessary condition for a minimiser of the energy functional. It is a PDE.
- ◆ The Euler-Lagrange equation for variational image restoration can be regarded as a fully implicit time discretisation of a nonlinear diffusion filter.
- ◆ The regularisation parameter  $\alpha$  approximates the stopping time of the diffusion filter.
- ◆ The regulariser  $\Psi(|\nabla u|^2)$  corresponds to the diffusivity  $\Psi'(|\nabla u|^2)$ .

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*(variational methods with a nonquadratic regulariser)*

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## Assignment T3 (1)



### Assignment T3 – Theoretical Home Work

#### Problem 1 (Discretisation of Anisotropic Nonlinear Diffusion)

(4 points)

For the discretisation of anisotropic nonlinear diffusion, one has to discretise the mixed derivative expression  $\partial_x(b\partial_y u) + \partial_y(b\partial_x u)$ .

Discretise the mixed derivative  $\partial_y(b\partial_x u)$  in two different ways:

- (a) ◆ using forward differences for  $\partial_y$  and  $\partial_x$ 
  - ◆ using backward differences for  $\partial_y$  and  $\partial_x$
  - ◆ averaging both results
- (b) ◆ using forward differences for  $\partial_y$  and backward differences for  $\partial_x$ 
  - ◆ using backward differences for  $\partial_y$  and forward differences for  $\partial_x$
  - ◆ averaging both results

Which approximations for the entire expression  $\partial_x(b\partial_y u) + \partial_y(b\partial_x u)$  are obtained with discretisations (a) and (b)? What is the result when these two approximations to  $\partial_x(b\partial_y u) + \partial_y(b\partial_x u)$  are averaged?

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## Assignment T3 (1)



### Problem 2 (Discretisation of Anisotropic Nonlinear Diffusion)

(2 points)

Find a suitable positive definite diffusion tensor  $\begin{pmatrix} a & b \\ b & c \end{pmatrix}$  where the “nonnegativity” semidiscretisation from Lecture 8 (slide 12) with  $h_1 = h_2$  has a negative stencil weight outside the centre. Can this also happen if the positive definite diffusion tensor is diagonally dominant?

### Problem 3 (Stopping Time Selection: Decorrelation Criterion)

(4 points)

Smooth the 1-D signal (3, 5, 3, 5, -5, -3, -5, -3) by discretised linear diffusion. Use the stencil (0.25, 0.5, 0.25) and reflecting boundary conditions. Apply the decorrelation criterion to determine when to stop the iterative process.

### Problem 4 (Convex Functionals and Forward Diffusion)

(2 points)

Consider the 1-D energy functional

$$E(u) = \int_a^b \left( (u - f)^2 + \alpha \Psi(u_x^2) \right) dx$$

and its corresponding diffusion process  $\partial_t u = \partial_x (\Psi'(u_x^2) u_x)$ . Assume that the regulariser  $s \mapsto \Psi(s^2)$  is convex in  $s$  such that  $\frac{d^2 \Psi}{ds^2} \geq 0$ . Show that the corresponding diffusion process always performs forward diffusion (and can therefore not act contrast-enhancing).

**Deadline for submission:** Friday, May 23, 10 am (before the lecture).

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